

Deformation of graded Poisson (Batalin- Vilkovisky) structures

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1. Introduction

We are interested in
the graded Poisson structures

Especially

the Batalin-Vilkovisky structures

Because

Mathematically

many geometrical and algebraic structures are described as a Batalin-Vilkovisky structure.

- Unify "all" geometries
- Classify "all" geometries
- Find a new geometry

Physically

a Batalin-Vilkovisky structure is the fundamental structure of gauge theories.

¹⁸¹ Batalin-Vilkovisky

- Unify "all" physics
- Classify "
- Find a new physics

Our tool is

deformation theory

based on the concept of
Kodaira - Spencer

First we consider
a simple BV-structure
(abelian topological sigma models)
and
construct all deformations.

Plan of my talk

§2 Graded Poisson Structures
on the Graded Vector Bundles

§3 Batalin-Vilkovisky Structures
of Abelian Topological Sigma Models

§4 Deformation Theory

§5 $n=2$

$n=3$

(§6 Quantum Version)

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2. Graded Poisson Structures on the Graded Vector Bundles (AKSZ formalism)

195 Alexandrov, Kontsevich Schwarz Zaboronky

M : a smooth manifold in d dimensions
 T^*M : a cotangent bundle

Def πT^*M is supermanifold
is a cotangent bundle with reversed parity of the fiber.

◦ local coordinate

$\{\phi^i\}$ on M $i = 1, 2, \dots, d$

$\{B_i\}$ on πT^*M

We assign a grading ^{the} **total degree**

$\{\phi^i\}$ is 0
 $\{B_i\}$ 1

More generally.

We define a graded cotangent bundle

$T^*[p]M$.

p : nonnegative integer.

the total degree of the fiber is p .

$\{\phi^i\}$ total degree 0

$\{B_i\}$ " p

Note) if p even $\Rightarrow B_i$: commutative

odd $\Rightarrow B_i$: anti commutative

• We can define

$T[p]M$: a graded tangent bundle

$E[p]$: a graded vector bundle

in a similar manner.

Generally we call a graded manifold

② P-structure

We consider a Poisson manifold N .

↓. shift the grading.
then We obtain a supermanifold \hat{N} and graded

$$N \longrightarrow \hat{N}$$

Poisson bracket

antibracket

$$\{*,*\}$$

$$\longrightarrow$$

$$(*,*)$$

It is called the P-structure.

Example 1 T^*M

T^*M has a natural Poisson (symplectic) structure:

$$\text{i.e. } \{F, G\} \equiv F \overset{\leftarrow}{\partial} \overset{\rightarrow}{\partial} G - F \overset{\leftarrow}{\partial} \overset{\rightarrow}{\partial} G$$

F, G : functions on T^*M

ϕ^i, B_i : Darboux coordinate

shift to

antibracket on $T^*[p]M$

P $(F, G) \equiv F \overset{\leftarrow}{\partial} \overset{\rightarrow}{\partial} G - F \overset{\leftarrow}{\partial} \overset{\rightarrow}{\partial} G$
 $\partial \phi^i \partial B_{p,i} \quad \partial B_{p,i} \partial \phi^i$

where $B_{p,i}$: total degree p

This antibracket satisfies

(*)
$$\begin{cases} (F, G) = -(-1)^{(|F|-p)(|G|-p)} (G, F) \\ (F, GH) = (F, G)H + (-1)^{(|F|-p)|G|} G(F, H) \\ (FG, H) = F(G, H) + (-1)^{|G|(|H|-p)} (F, H)G \\ (-1)^{(|H|-p)(|H|-p)} (F, (G, H)) \\ + (\text{cyclic permutations of } FGH) = 0 \end{cases}$$

where $(|F|, |G|, |H|)$ are the total degrees of F, G & H .

(*,*) carries the total degree $-p$

Def An antibracket on a graded mfd \tilde{N} with a total degree $-p$ is a bilinear form which satisfies (*).

Example 2 $E \oplus E^*$

has a natural pairing on the fiber

$\{A^a\}$ local coordinate on the fiber of E

$\{B_a\}$ local coordinate on the fiber of E^*

We can define a Poisson bracket

$$\{F, G\} \equiv F \overset{\leftarrow}{\partial} \overset{\rightarrow}{\partial} G - F \overset{\leftarrow}{\partial} \overset{\rightarrow}{\partial} G$$

$\frac{\partial}{\partial A^a} \quad \frac{\partial}{\partial B_a} \quad \frac{\partial}{\partial B_a} \quad \frac{\partial}{\partial A^a}$

If We consider a graded bundle

$$E[p] \oplus E^*[q]$$

the Poisson bracket shifts to an antibracket

$$(F, G) \equiv F \overset{\leftarrow}{\partial} \overset{\rightarrow}{\partial} G - (-1)^{p_F} F \overset{\leftarrow}{\partial} \overset{\rightarrow}{\partial} G$$

$\frac{\partial}{\partial A_p^a} \quad \frac{\partial}{\partial B_{q,a}} \quad \frac{\partial}{\partial B_{q,a}} \quad \frac{\partial}{\partial A_p^a}$

$(*, *)$ total degree $-p-q$

$a = 1, \dots, \text{rank } E$

this antibracket satisfies (\star) .

② Q-structure

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Def S : a function on a graded bundle is called a Batalin-Vilkovisky (BV) action (an action) if

S satisfies
the classical master equation
 $(S, S) = 0$

It is called Q-structure.

Def a Batalin-Vilkovisky structure is P-structure & Q-structure on a graded manifold \tilde{N} .

Def $(S, *) = \delta(*)$ is called the BRST transformation which defines coboundary operator δ s.t. $\delta^2 = 0$ $\therefore (S, S) = 0$

3 Batalin-Vilkovisky Structures of Abelian Topological Sigma Models ⌋

Now we consider a sigma model.

We introduce two ^(topological) manifolds

X and M .

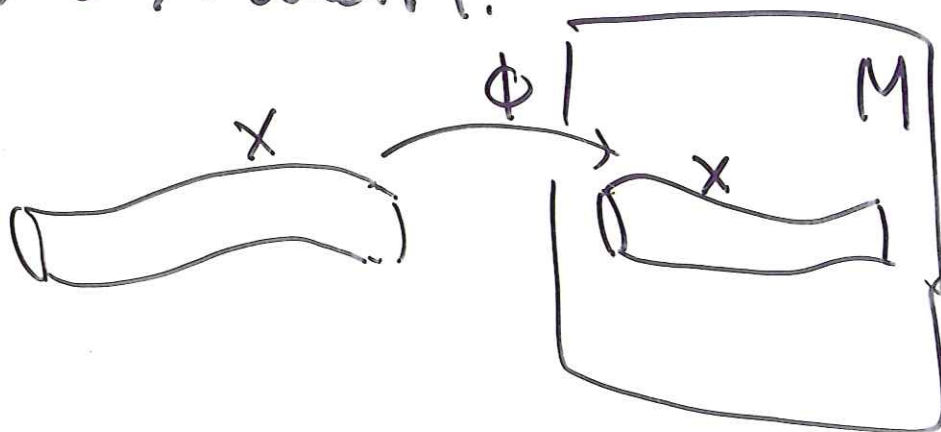
And consider a smooth map $\phi: X \rightarrow M$.

o A sigma model

is a classical or quantum field theory constructed from a map ϕ on X (& auxially other maps).

o A topological sigma model

is a sigma model independent of metrics on X and M .



X is n -dim is called worldsheet ($n=2$)⁽¹²⁾
world volume ($n \geq 3$)

M is d -dim target space

① We extend a map $\phi: X \rightarrow M$
to $\Phi: \pi TX \rightarrow M$

Next
② We construct the BV structure
from $\Phi: TX \rightarrow M$

From now on, we consider
two cases T^*M & $E \oplus E^*$

3.1 P-structure

Example 1 T^*M

The dimension n of X labels a natural grading on $T^*[P]M$ as $p = n - 1$

Antibracket on $T^*[n-1]M$

$$(F, G) \equiv F \overset{\leftarrow}{\partial} \overset{\rightarrow}{\partial} G - F \overset{\leftarrow}{\partial} \overset{\rightarrow}{\partial} G$$

$$\frac{\partial}{\partial \Phi^i} \frac{\partial}{\partial B_{n,i}} - \frac{\partial}{\partial B_{n,i}} \frac{\partial}{\partial \Phi^i}$$

$$\Phi^i \in \pi T^*X \otimes M$$

$$B_{n,i} \in \pi \left(\wedge^{n-1} \pi T^*X \otimes \Phi^*(T^*[n-1]M) \right)$$

(,) has the total degree $-n+1$

Example 2 $E[p] \oplus E^*[q]$

In order to make the total degree of an antibracket $-n+1$, we take

$$q = n - p - 1.$$

We consider

$$E[p] \oplus E^*[n-p-1]$$

P the antibracket is

$$(F, G) \equiv F \overset{\leftarrow}{\partial} \frac{\overset{\rightarrow}{\partial}}{\partial A_p^{ap}} G - (-)^{np} F \overset{\leftarrow}{\partial} \frac{\overset{\rightarrow}{\partial}}{\partial B_{n-p-1, ap}} G$$

where $A_p^{ap} \in \Gamma(\wedge^p T^*X \otimes \phi^*(E[p]))$

$B_{n-p-1, ap} \in \Gamma(\wedge^{n-p-1} T^*X \otimes \phi^*(E^*[n-p-1]))$

$(,)$ has the total degree $-n+1$.

We want to consider $T^*[n-1]M$
 & $E[p] \oplus E^*[n-p-1]$ simultaneously.
 ($p=1, 2, \dots$)

We identify

$$E^*[n-p-1] \oplus (E^*)^*[p] \cong E[p] \oplus E^*[n-p-1]$$

as a dual space.

$\lfloor x \rfloor$: the floor fn which gives the largest integer less than or equal to x .

Let E_p be $\lfloor \frac{n-1}{2} \rfloor$ series of vector bundles, $1 \leq p \leq \lfloor \frac{n-1}{2} \rfloor$

We consider a vector bundle

$$\left(\sum_{p=1}^{\lfloor \frac{n-1}{2} \rfloor} E_p[p] \oplus E_p^*[n-p-1] \right) \oplus T^*[n-1]M$$

the combination of Ex, 1 & 2

The antibracket is the combination of Ex. 1 & 2

$$\boxed{P} \quad (F, G) \equiv \sum_{p=0}^{\lfloor \frac{n-1}{2} \rfloor} F \overset{\leftarrow}{\partial} \overset{\rightarrow}{\partial} G - (-1)^{np} F \overset{\leftarrow}{\partial} \overset{\rightarrow}{\partial} G$$

$$\frac{\partial}{\partial A_p^{a_p}} \frac{\partial}{\partial B_{np+1, a_p}} \quad \frac{\partial}{\partial B_{np+1, a_p}} \frac{\partial}{\partial A_p^{a_p}}$$

where $p=0$ part is the antibracket on

$$T^*([n-1]M, \mathbb{A}_0^{a_0} = \phi^i$$

(,) total degree $-n+1$

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Def the **form degree** is a grading
on ΠT^*X

the form degree 0 on the base mfd X
" 1 on the fiber $\Pi T^*_\sigma X$

$\deg \bar{F}$: form degree of \bar{F}

Def the **ghost number** $gh F$ is
defined $gh F = |F| - \deg \bar{F}$

3-2 \mathbb{Q} -structure

A simplest natural \mathbb{Q} -structure is

$$S_0 = \sum_{p=0}^{\lfloor \frac{n-1}{2} \rfloor} (-1)^{n-p} \int_X B_{n-p, p} dA_p^{ap}$$

where d : exterior differential on X

\int_X : integration over the n -form part

analogy of $\theta = p_i dq^i$
 a fundamental form for a symplectic form
 $\omega = dq^i \wedge dp_i = -d\theta$.

It is called an abelian BF theory in n dim

We can conform

$$(S_0, S_0) = 0$$

i.e. S_0 is \mathbb{Q} -structure.

$|S_0| = n$ total degree

4. Deformation Theory

'99 Izawa
'00~ N.I.

Def Deformation of the BV structures

Fix a P-structure $(*, *)$ on a graded bundle.

Consider infinitesimal deformations of Q-structure where $S = \int_X F(A_p, B_{q-p})$

$$S = S_0 + gS_1 + g^2S_2 + \dots$$

such that $(S, S) = 0$

$$\begin{aligned} \text{up to } S' &= S + \delta_0(*) \\ \Rightarrow S' &\sim S \quad \delta_0 \text{ exact} \\ &\text{equivalent} \end{aligned}$$

where g : constant (deformation parameter)

$$\delta_0(*) \equiv (S_0, *)$$

Note) $\delta_0^2 = 0$; We look for δ_0 -cohomology class.

$$S = S_0 + gS_1 + g^2 S_2 + \dots$$

We set $g S_i'$

We consider the total bundle

$$\sum_{p=1}^{\lfloor \frac{n-1}{2} \rfloor} (E_p[p] \oplus E_p^*[n-p]) \oplus T^*[n-1]M$$

~~P-structure~~

$$(F, G) = \sum_{p=0}^{\lfloor \frac{n-1}{2} \rfloor} F \begin{array}{c} \leftarrow \quad \rightarrow \\ \frac{\partial}{\partial A_p^{ap}} \quad \frac{\partial}{\partial B_{n-p, ap}} \end{array} G \quad (-) \quad F \begin{array}{c} \leftarrow \quad \rightarrow \\ \frac{\partial}{\partial B_{n-p, ap}} \quad \frac{\partial}{\partial A_p^{ap}} \end{array} G$$

Q-structure

$$S_0 = \sum_{p=0}^{\lfloor \frac{n-1}{2} \rfloor} (-1)^{np} \int_X B_{n-p, ap} dA_p^{ap}$$

Now we assume

2!

1. For an arbitrary function G

$$\int_X dG(\Phi, d\Phi) = 0$$

where $\Phi = A_p^{a_p}$ or B_{n-p-1, a_p}

no obstruction

$$\Leftrightarrow H^2(\delta_0) = 0 \quad \text{no D-brane}$$

2. total degree

$$|S| = |S_1| = n$$

is physically required but

mathematically can be relaxed.

Procedure

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$$S = S_0 + gS_1'$$

Since

$$(S, S) = 0 \Rightarrow (S_0 + gS_1', S_0 + gS_1') = 0$$

We obtain

$$(i) \underline{g^0} \quad (S_0, S_0) = 0 \quad \underline{OK}$$

$$(ii) \underline{g^1} \quad (S_0, S_1') = 0$$

Since

$$(S_0, A_p^{ap}) = dA_p^{ap}$$

$$(S_0, B_{n-p-1, ap}) = dB_{n-p-1, ap}$$

The general solution is

$$S_1' = \int_X F(\Phi, d\Phi)$$

F: an arbitrary function

$$(S_0, S_1') = 0$$

Assumption 1

Theorem

a monomial $F(\Phi, d\Phi)$ includes at least one $d\Phi$ and $(S_0, \int_X F) = 0$,
 $\int_X F(\Phi, d\Phi)$ is δ_0 -exact.

\therefore F includes at least one $d \Rightarrow$

$$\int_X F(\Phi, d\Phi) = \sum_{p=0}^{n-1} \int_X F_{n-p-1} dG_p$$

F_{n-p-1} : form degree $n-p-1$

G_p : " " p

expansion by the form degrees

$$(S_0, \int_X F) = 0$$

$$\Rightarrow \left\{ \begin{array}{l} \delta_0 F_0 = 0 \\ \delta_0 F_{n-p-1} = d F_{n-p-2} \quad -1 \leq p \leq n-2 \\ d F_n = 0, \delta_0 G_0 = 0 \\ \delta_0 G_p = d G_{p-1} \quad 1 \leq p \leq n \\ d G_n = 0 \end{array} \right.$$

for even p adjacent 2 terms

(29)

$$\begin{aligned} & F_{n-p-1} dG_p + F_{n-p-2} dG_{p+1} \\ (*) \\ & = (-1)^{n-p-1} \delta_0(F_{n-p-1} G_{p+1}) - (-1)^{n-p} d(F_{n-p-2} G_{p+1}) \end{aligned}$$

Integration of the above is δ_0 -exact

$$\text{If } n \text{ even} \Rightarrow \sum_{p=0}^{n-1} F_{n-p-1} dG_p$$

even numbers of terms

\Rightarrow Integration is δ_0 -exact

$$\text{If } n \text{ odd} \Rightarrow \sum_{p=0}^{n-1} F_{n-p-1} dG_p$$

odd number of terms

but the last term

$$F_0 dG_n \stackrel{(*)}{=} \delta_0(F_0 G_n)$$

\Rightarrow Integration is δ_0 -exact //

From the theorem, the general solution (25)

$$S_1' = \int_x F(\Phi)$$

F : an arbitrary function

$$S_1' = \sum \int_x F_{a_{p(1)} \dots a_{p(k)}}^{b_{z(1)} \dots b_{z(k)}} (A_0^{a_p})$$

$$\times A_{p_1}^{a_{p_1}} \dots A_{p_k}^{a_{p_k}} (B_{z_1, b_{z_1}} \dots B_{z_k, b_{z_k}})$$

the total degree n terms

(assumption 2)

arbitrary fn F for $A_0^{a_p} = \Phi^{a_p}$

$$(ii) \quad g^2 (S_1', S_1') = 0 \quad (26)$$

\Rightarrow We obtain the conditions for the functions $F_{a \dots a}^{b \dots b} (A_0^{ap})$, which determine the algebraic and geometric structure.

The solution

$$S = S_0 + g S_1' \\ = \sum_{p=0}^{\lfloor \frac{n-1}{2} \rfloor} (-1)^{u-p} \int_X B_{n-p, ap} dA_p^{ap}$$

$$+ g \int_X F(\Phi)$$

$$\& (S_1', S_1') = 0$$

5 Structures in lower n

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5-1 $n=2$

the total bundle

$$T^*[1]M = \Pi T^*M$$

P

$$(F, G) = F \frac{\partial}{\partial \Phi^i} \frac{\partial}{\partial B_{ij}} G - F \frac{\partial}{\partial B_{ij}} \frac{\partial}{\partial \Phi^i} G$$

Q

$$S = S + g S'$$

$$= \int_x B_{ij} d\Phi^i + g \int_x \frac{1}{2} f^{ij}(\Phi) B_{ij}$$

$$(S', S') = 0$$

$$\Leftrightarrow f^{kl} \frac{\partial}{\partial \Phi^l} f^{ij} + f^{il} \frac{\partial}{\partial \Phi^l} f^{jk} + f^{jl} \frac{\partial}{\partial \Phi^l} f^{ki} = 0$$

$$\Leftrightarrow -f^{ij} \frac{\partial}{\partial \Phi^i} \frac{\partial}{\partial \Phi^j}$$

is a Poisson bivector

the Poisson sigma model

$$\{F, G\} = -f^{ij}(\Phi) \frac{\partial F}{\partial \Phi^i} \frac{\partial G}{\partial \Phi^j} = ((S, F), G)$$

f. Lie algebroid and observables

Def Lie algebroid

- vector bundle $\mathcal{E} \rightarrow M$
- for section $e_1, e_2 \in \Gamma(\mathcal{E})$
the bracket $[e_1, e_2]$ is defined
with a Lie algebra str.

- bundle map (anchor)

$$\rho: \mathcal{E} \rightarrow TM$$

(1) s.t. for $\forall e_1, e_2 \in \Gamma(\mathcal{E})$

$$[\rho(e_1), \rho(e_2)] = \rho([e_1, e_2])$$

2) for $\forall e_1, e_2 \in \Gamma(\mathcal{E}), F \in C^\infty(M)$

$$[e_1, Fe_2] = F[e_1, e_2] + (\rho(e_1)F)e_2$$

If we define

$$[e_1, e_2] = ((S, e_1), e_2)$$

$$f(e) F(\phi) = (e, (S, F(\phi)))$$

where S : Poisson sigma model

\mathcal{E} is Lie algebroid

$$\iff (S, S) = 0$$

where $M = \{ \phi^i : \Sigma \rightarrow M \}$

$$\mathcal{E} = T^*M$$

Batalin-Vilkovisky structure
of total degree 2 topological
 σ -model

\cong Lie algebroid on T^*M

Levin, Olshansky

5-2, n=3

the total bundle

$$(E[1] \oplus E^*[1]) \oplus T^*[2]M$$

P

$$\begin{aligned}
(F, G) = & F \overset{\leftarrow}{\partial} \overset{\rightarrow}{\partial} G - F \overset{\leftarrow}{\partial} \overset{\rightarrow}{\partial} G \\
& + F \overset{\leftarrow}{\partial} \overset{\rightarrow}{\partial} G + F \overset{\leftarrow}{\partial} \overset{\rightarrow}{\partial} G
\end{aligned}$$

$\frac{\partial \Phi^i}{\partial B_{2i}} \frac{\partial}{\partial B_{2i}} G$ $\frac{\partial}{\partial B_{2i}} \frac{\partial \Phi^i}{\partial B_{2i}} G$
 $\frac{\partial}{\partial A_i} \frac{\partial}{\partial B_{1i}} G$ $\frac{\partial}{\partial B_{1i}} \frac{\partial}{\partial A_i} G$

Q

$$S = S_0 + gS_1'$$

$$S_0 = \int_x -B_{2i} d\Phi^i + B_{1a} dA_1^a$$

$$\begin{aligned}
S_1' = & \int f_{1a}^i(\Phi) A_1^a B_{2i} + f_2^{ib}(\Phi) B_{2i} B_{1b} \\
& + \frac{1}{3!} f_{3abc}(\Phi) A_1^a A_1^b A_1^c + \frac{1}{2} f_{4ab}^c(\Phi) A_1^a A_1^b B_{1c} \\
& + \frac{1}{2} f_{5a}^{bc}(\Phi) A_1^a B_{1b} B_{1c} + \frac{1}{3!} f_6^{abc}(\Phi) B_{1a} B_{1b} B_{1c}
\end{aligned}$$

$$(S'_i, S'_i) = 0 \Leftrightarrow$$

$$\textcircled{1} f_{1e}^i f_2^j e + f_2^{ie} f_{1e}^j = 0$$

$$\textcircled{2} - \frac{\partial f_{1c}^i}{\partial \phi^j} f_{1b}^j + \frac{\partial f_{1b}^i}{\partial \phi^j} f_{1c}^j + f_{1e}^i f_{4bc}^e + f_2^{ie} f_{3ebc} = 0$$

$$\textcircled{3} f_{1b}^j \frac{\partial f_2^{ic}}{\partial \phi^j} - f_2^{jc} \frac{\partial f_{1b}^i}{\partial \phi^j} + f_{1e}^i f_{5b}^{ec} - f_2^{ie} f_{4eb}^c = 0$$

$$\textcircled{4} - f_2^{jb} \frac{\partial f_2^{ic}}{\partial \phi^j} + f_2^{jc} \frac{\partial f_2^{ib}}{\partial \phi^j} + f_{1e}^i f_6^{ebc} + f_2^{ie} f_{5e}^{bc} = 0$$

$$\textcircled{5} - f_{1ca}^j \frac{\partial f_{4bc}^d}{\partial \phi^j} + f_2^{jd} \frac{\partial f_{3abc}}{\partial \phi^j} + f_{4ea}^d f_{4bc}^e + f_{3cab} f_{5c}^{de} = 0$$

$$\textcircled{6} - f_{1ca}^j \frac{\partial f_{5b}^{cd}}{\partial \phi^j} + f_2^{jc} \frac{\partial f_{4ab}^d}{\partial \phi^j} + f_{3cab} f_6^{ecd} + f_{4ea}^c f_{5b}^{cde} + f_{4ab}^e f_{5e}^{cd} = 0$$

$$\textcircled{7} - f_{1a}^j \frac{\partial f_6^{bcd}}{\partial \phi^j} + f_2^{jb} \frac{\partial f_{5a}^{cd}}{\partial \phi^j} + f_{4ea}^c f_6^{bcd} + f_{5e}^{bc} f_{5a}^{cde} = 0$$

$$\textcircled{8} - f_2^{ja} \frac{\partial f_6^{bcd}}{\partial \phi^j} + f_6^{eab} f_{5e}^{cd} = 0$$

$$\textcircled{9} - f_{1ca}^j \frac{\partial f_{3bcd}}{\partial \phi^j} + f_{4ab}^e f_{3cd}^e = 0$$

Theorem

$$(S_1', S_1') = 0 \Leftrightarrow \textcircled{1} \sim \textcircled{9}$$

\Leftrightarrow a Courant algebroid
on $E[1] \oplus E^*[1]$

$\Leftrightarrow n=3$ Batalin-Vilkovisky structure

Courant algebroid

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197 Liu, Weinstein, Xu
01 Roytenberg 33

vector bundle $E \rightarrow M$

with $\langle \cdot, \cdot \rangle$: (graded) symmetric bilinear form

\circ : bilinear form (Dorfman bracket)

ρ : $E \rightarrow TM$ anchor

s.t.

$$1. e_1 \circ (e_2 \circ e_3) = (e_1 \circ e_2) \circ e_3 + e_2 \circ (e_1 \circ e_3)$$

$$2. \rho(e_1 \circ e_2) = [\rho(e_1), \rho(e_2)]$$

$$3. e_1 \circ F e_2 = F(e_1 \circ e_2) + (\rho(e_1) F) e_2$$

$$4. e_1 \circ e_2 = \frac{1}{2} \mathcal{D} \langle e_1, e_2 \rangle$$

$$5. \rho(e_1) \langle e_2, e_3 \rangle = \langle e_1 \circ e_2, e_3 \rangle + \langle e_2, e_1 \circ e_3 \rangle$$

where $e_1, e_2, e_3 \in \Gamma(E)$

$$F \in C^\infty(M)$$

$$\mathcal{D} : M \rightarrow \Gamma(E) \text{ s.t. } \langle \mathcal{D}F, e \rangle = \rho(e)F$$

We take

$$M = \{ \phi : X \rightarrow M \}$$

fiber of \mathcal{E} : $E[1] \oplus E^*[1]$

basis $\{A^a, B^a\}$

If we set

$$\left\{ \begin{aligned} \langle e_1, e_2 \rangle &\equiv (e_1, e_2) && \text{antibracket} \\ e_1 \circ e_2 &\equiv ((S, e_1), e_2) \\ \rho(e) F(\phi) &\equiv (e, (S, F(\phi))) \\ \mathcal{D}(\ast) &\equiv (S, \ast) \end{aligned} \right\} \text{ defined as derived bracket}$$

then

$$\langle, \rangle, \circ, \rho, \mathcal{D} \text{ is Courant algebroid}$$

$$\iff (S, S) = 0 \quad \text{master eq}$$

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Therefore S is $|S|=3$

We call

S : Courant sigma model

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The form degree n deformation defines an algebroid on the bundle defined from the master eq. $\text{on } M$

$$(S, S) = 0$$

We name BV- n -algebroid

n -algebroid	classical	quantum
2 Lie algebroid	Poisson	*-product on Poisson
3 Courant	Courant algebroid	?
n n -algebroid	BV- n -algebroid	?

② Sub structures of BV structures
of symplectic (Kähler)

A-model

$$S = \int_{X_2} B_{ij} d\phi^i + \int_{X_2} \frac{1}{2} f^{ij} B_{ij}$$

o complex

B-model $i = (a, \bar{a})$

$$S = \int_{X_2} B_{ia} d\phi^a + B_{i\bar{a}} d\phi^{\bar{a}} - B_{0\bar{a}} dA_i^{\bar{a}} - B_{i\bar{a}} A_i^{\bar{a}}$$

o twisted Poisson

$$S = \int_{X_2} B_{ij} d\phi^i + \frac{g}{2} f^{ij} B_{ij} + \int_{X_3} \frac{1}{2} H_{ijk} d\phi^i d\phi^j d\phi^k$$

$n=3$

generalized complex structure

$$\begin{aligned}
S = & \int_{X_3} -B_{zi} d\phi^i + B_{i\bar{c}} dA_1^{\bar{c}} \\
& + \int_{X_3} -J_{\bar{j}}^i B_{zi} A^{\bar{j}} - P^{\bar{i}\bar{j}} B_{zi} B_{i\bar{j}} \\
& + \frac{1}{2} \left(H_{ijk} + \frac{\partial Q_{jk}}{\partial \phi^i} \right) A_1^i A_1^{\bar{j}} A_1^k \\
& + \frac{1}{2} \left(-\frac{\partial J_{\bar{j}}^i}{\partial \phi^i} + \frac{\partial J_{\bar{i}}^k}{\partial \phi^{\bar{j}}} \right) A_1^i A_1^{\bar{j}} B_{ik} \\
& + \frac{1}{2} \frac{\partial P_{\bar{i}k}}{\partial \phi^{\bar{i}}} A_1^i B_{i\bar{j}} B_{ik}
\end{aligned}$$

where $J = \begin{pmatrix} J & P \\ Q & -J^* \end{pmatrix}$ GCS

etc.

6 Quantum Deformation Theory 138

Def BV Laplacian

$$\Delta F \equiv \sum_{p=0}^{n-1} \frac{\overrightarrow{\partial}}{\partial A_p} \overleftarrow{\partial}_{B_{n-p}}$$

property

$$\Delta(FG) = (\Delta F)G + (-1)^{(n+1)F} (FG) + (-1)^F F\Delta G$$

We add the gauge fixing term S_{GF}

$$S_g = S + S_{GF}$$

to restrict the theory on the "Lagrangian" submanifold of the total bundle.

S_g satisfies **the quantum master equation**

$$(S_g, S_g) - 2i\hbar \Delta S_g = 0$$

$$\hbar = g$$

Def observable \mathcal{O} : fn on the total bundle

$$(S_g, \mathcal{O}) - i\hbar \Delta \mathcal{O} = 0$$

the "correlation function"

$$\langle \mathcal{O}_1, \dots, \mathcal{O}_k \rangle \quad \text{path integral}$$

$$= \int \prod_{p=0}^{\lfloor \frac{n-1}{2} \rfloor} \mathcal{D}A_p \mathcal{D}B_{n-p-1} \mathcal{O}_1 \dots \mathcal{O}_k e^{\frac{i}{\hbar} S_g}$$

define the geometric or algebraic str.

@ $n=2$ known

quantization of the Poisson sigma model
on $X = \text{disc}$

\Leftrightarrow deformation quantization
on a Poisson manifold

'97 Kontsevich
'99 Cattaneo-Felder

In $n=2$

$$S_0$$

↓ deformation

$$S = S_0 + g S_1$$

provides a quantum deformation

$$F(x) G(x) = \int_{\phi(x)=x} \mathcal{D}\phi \mathcal{D}\psi \mathcal{D}B, F(\phi(x)) G(\phi(x)) \times e^{\frac{i}{\hbar} (S_0 + S_{GF})}$$

↓ star deformation

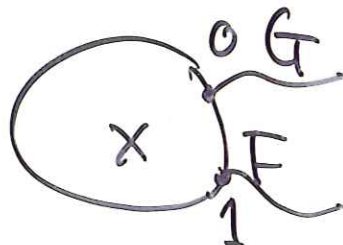
$$F \star G(x) = \int_{\phi(x)=x} \mathcal{D}\phi \mathcal{D}\psi \mathcal{D}B, F(\phi(x)) G(\phi(x)) \times e^{\frac{i}{\hbar} (S + S_{GF})}$$

$$\phi(x) = x$$

F, G : fn on M

Observables

$$\hbar = g$$



8 Other Results

- Dimensional reduction $n=3 \rightarrow n=2$

Covariant algebraic is formulated as a $n=2$ topological sigma model

- complex structure (B-model)
- Kähler structure (A-model)
- generalized geometry (Kapstan, Zuckani, Nil.)
- twisted Poisson str. (WZ-Poisson)
- contact (Jacobi) str. (Jacobi σ -model)

are formulated as a sub structure of our model

- Many geometric structures can be formulated as a BV structure

\Rightarrow unification of geometries?
(theory of every geometries?)

n	classical	quantum
2	Poisson Complex B Kähler A Contact (Jacobi) Symplectic -----	*-product B-model A-model
3	Courant Dirac generalized geometry twisted Poisson topological M G_2 -----	<div style="border: 1px solid black; border-radius: 50%; padding: 10px; display: inline-block;"> Corresponding quantum Geometry </div>
n	topological membrane	→

The list includes many geometries but have not completed yet.

Problems

- Characterization of n -algebroid
classically ($n \geq 4$)
quantum ($n \geq 3$)
- Extension to the other graded bundles.

Formulate a given geometry
as a (graded) Poisson geometry.